

# $H_2$ molecules and cold clouds in cooling flow clusters<sup>†</sup>

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## 1. Introduction

Around the epoch of recombination atomic hydrogen is the most important chemical species and leads by adiabatic cooling of the universe to the formation of molecular hydrogen  $H_2$  (Puy et al. 1993, Puy & Signore 1999). The actual  $H_2$  content is very uncertain and estimated only indirectly. The important recent observation of the lowest pure rotational lines of  $H_2$  in the spiral galaxy NGC 891 (Valentijn & Van der Werf 1999) gives a direct indication of relatively warm ( $T=150-230$  K) molecular clouds in the disk in addition to a massive cooler (80-90 K) component in the outer regions.

In clusters of galaxies X-ray measurements show an excess absorption below  $\sim 1$  keV compared to a best fit bremsstrahlung model, which is interpreted as due to the presence of cold molecular clouds (White et al. 1991).

In a scenario with successive fragmentation of these clouds we calculate the molecular rotational line cooling due to  $HD$ - and  $H_2$ -molecules and determine their minimum temperature achievable in equilibrium with the exterior bremsstrahlung of the hot intracluster gas.

## 2. Molecular cooling

The abundance  $\eta_{HD}$  of the  $HD$  molecule is considered to be primordial:

$$\eta_{HD} \simeq 7 \times 10^{-5}, \text{ Signore \& Puy 1999.}$$

The quadrupole transitions of  $H_2$  are supposed to follow an ortho/para ratio of 1, its first excited state is at 512 K. Although this cooling is less efficient than the cooling due to  $HD$ ,  $H_2$  is more abundant and we expect the  $H_2$  cooling to be more important than the  $HD$  cooling in the temperature region above  $\sim 100$  K and in the density region of  $10^4 \text{ cm}^{-3}$ .

Considering only the transition between the ground state and the first rotational level, the molecular cooling can be calculated analytically (Puy, Grenacher and Jetzer 1999). In the case of  $H_2$ -clouds temperatures as mentioned above, the higher excitation levels turn out to be significant even if they are still very weakly populated.

In Figure 1 we show for the total cooling  $\Lambda_{H_2} + \Lambda_{HD}$  the relative importance of the cooling agents  $H_2$  and  $HD$

$$\alpha_{H_2} = \frac{\Lambda_{H_2}}{\Lambda_{H_2} + \Lambda_{HD}}$$

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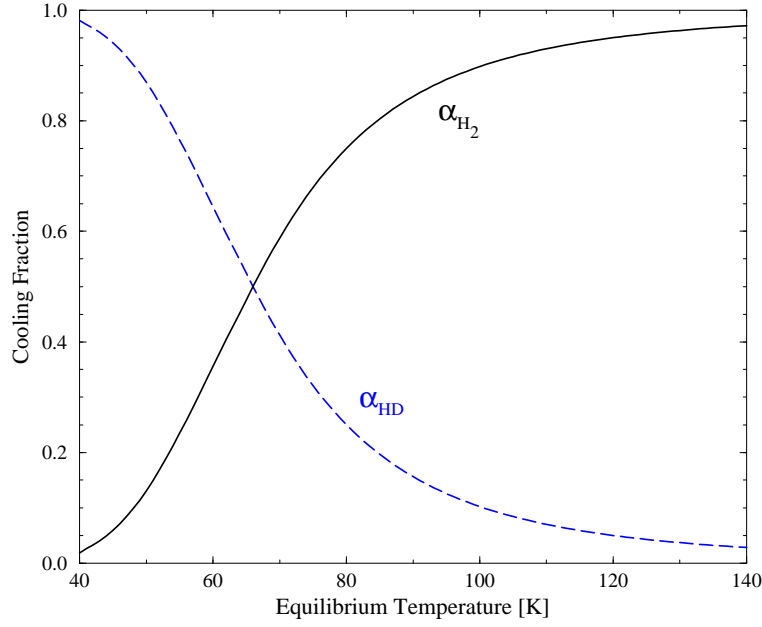


FIGURE 1. Relative importance of the different cooling agents  $H_2$ ,  $\alpha_{H_2} = \Lambda_{H_2}/(\Lambda_{H_2} + \Lambda_{HD})$  and  $HD$ ,  $\alpha_{HD} = \Lambda_{HD}/(\Lambda_{H_2} + \Lambda_{HD})$ .

$$\alpha_{HD} = \frac{\Lambda_{HD}}{\Lambda_{H_2} + \Lambda_{HD}} .$$

We can see clearly, that  $H_2$  becomes the most important cooling agent above  $\sim 70$  K.

### 3. Thermal equilibrium

The most important heat source for molecular clouds in cooling flows is the X-ray bremsstrahlung emitted by the hot intracluster gas. In the successive fragmentation scenario this radiation is shielded by the presence of an attenuating column density in the outer parts of the clouds, which we take into account by the attenuation factor,  $\tau$ , following O’Dea et al. (1994).

The balance between heating and cooling in the cluster environment leads to a thermal equilibrium inside the cooling flow region of the clusters. The escape probability for this molecular emission is close to 1. We investigate the coldest equilibrium achievable inside the cooling flow region.

### 4. Discussion

We consider small clouds ( $\sim 10$  AU) with an  $H_2$  density of  $10^4 \text{ cm}^{-3}$ , attenuated by  $\tau = 0.01$ . The X-ray heating ( $\Gamma_X(r) \propto r^{-3}$ ) is of course more important in the cluster center and low equilibrium temperatures are achieved at large distances from the cluster center. Nevertheless, the fact that the cold clouds are located in the cooling flow region implies that the distance must be below the cooling radius. In this context we have calculated the equilibrium temperature of the molecular clouds located at the cooling radius. The table gives these equilibrium temperatures  $T_{clump}$  for different clusters of galaxies.

Cluster	$T_{\text{clump}}$ (in K)	$r_{\text{cool}}$ (in kpc)	$T_{\text{Kev}}$ (in keV)
Centaurus	153	87	2.1
Hydra A	107	162	4.5
PKS 0745-191	92	214	8.6
Abell 262	119	67	2.5
Abell 426	83	145	6.3
Abell 478	228	240	7.1
Abell 496	58	138	4.8
Abell 539	133	34	3.4
Abell 576	102	69	2.9
Abell 1060	102	68	3.3
Abell 1367	80	40	4.1
Abell 1795	136	181	5.3
Abell 2052	96	140	3.4
Abell 2151	256	146	2.9
Abell 2159	128	119	4.5

TABLE 1. The equilibrium temperature  $T_{\text{clump}}$  at the cooling radius  $r_{\text{cool}}$  for the cluster temperature  $T_{\text{keV}}$  are shown for different cluster of galaxies.

In the region of the cooling flow, we find that an equilibrium is possible at low temperature (below 200 K) due to  $H_2$  cooling.

Whether the cloud can be cooled down to  $\sim 70$  K, where  $HD$  dominates, depends on various parameters, such as the cluster temperature  $T_{\text{keV}}$ , the attenuation factor  $\tau$ , the cooling radius  $r_{\text{cool}}$  and other characteristics of the hot intracluster gas.

The detection of these  $H_2$  molecules is difficult, because the first rotational level, accessible only through a quadrupolar transition, is more than 500 K above the fundamental. The study of  $CO$ - $H_2$  ratios can give some insight, because the  $CO$  molecules are excited by collisions with  $H_2$ , and should be a tracer of cold  $H_2$  clouds (Grenacher et al. 1999). In this context cold  $H_2$ -clouds could be an interesting possibility of baryonic dark matter (Combes 1999). The FUSE satellite will certainly give clarification on this problem.

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